



# Controllable low-temperature growth and enhanced photoelectrochemical water splitting of vertical SnS<sub>2</sub> nanosheets on graphene

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## ABSTRACT

Two-dimensional (2D) SnS<sub>2</sub> nanosheets vertically grown on graphene electrodes have still not been demonstrated despite their considerable research interest and potential for photoelectrochemical (PEC) applications. Herein, the controllable growth and enhanced PEC water-splitting performance of vertically aligned 2D SnS<sub>2</sub> nanosheets on graphene are reported. The heterojunction quality of the graphene/2D SnS<sub>2</sub> nanosheets was ensured by low-temperature growth at 230 °C using metalorganic chemical vapor deposition, resulting in significantly improved charge transfer properties. The PEC photocurrent density and photoconversion efficiency significantly increased upon insertion of a graphene layer. In addition, the 2D SnS<sub>2</sub> nanosheet microscopic structure dependency of the PEC performance was systematically studied. The best PEC reactivity of the 2D SnS<sub>2</sub> nanosheets was achieved at a height of ~0.85 μm, corresponding to nearly the limit of the carrier diffusion length, and a thickness of ~70 nm to balance the low interfacial contact resistance with an effective photogenerated carrier dynamics. These results, including the low-temperature vertical growth of 2D SnS<sub>2</sub> on graphene, can effectively be utilized to exploit the full potential of 2D SnS<sub>2</sub> for various PEC applications.

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## 1. Introduction

Photoelectrochemical (PEC) water splitting has attracted considerable attention as a promising eco-friendly technology for converting solar energy to clean hydrogen fuel [1–3]. Numerous semiconductors, including oxides and metal chalcogenides, have been extensively studied for PEC photoelectrode applications [3]. Recently, graphene-like two-dimensional (2D) metal dichalcogenides, such as SnS<sub>2</sub>, MoS<sub>2</sub>, and WSe<sub>2</sub>, have emerged as promising semiconductor photocatalysts owing to their strong photoexcitation, good chemical stability, and earth abundance [4–7]. 2D SnS<sub>2</sub> is of particular interest due to its appropriate bandgap energy of 2.18–2.44 eV, high carrier mobility (~300 cm<sup>2</sup> V<sup>−1</sup> s<sup>−1</sup> for 3-nm-thick SnS<sub>2</sub>), high optical absorption, and strong photocatalytic activity [7,8]. SnS<sub>2</sub>, an n-type semiconductor, is generally utilized as a photoanode, at which an oxygen evolution reaction (OER: H<sub>2</sub>O → 2H<sup>+</sup> + ½O<sub>2</sub> + 2e<sup>−</sup>) occurs, whereas a hydrogen evolution reaction (HER: 2H<sup>+</sup> + 2e<sup>−</sup> → H<sub>2</sub>) occurs at the counter cathode, e.g., Pt. The appropriate design of 2D SnS<sub>2</sub> nanosheets on conductive substrates in terms of physical nanostructure architecture and the

energy band diagram is the decisive factor for improving photo-generation and transportation of charge carriers to electrode/liquid interfaces, consequently enhancing the OER and HER.

The ideal architecture of 2D SnS<sub>2</sub> nanosheets on conductive substrates involves vertically standing 2D SnS<sub>2</sub> structures. The vertically standing 2D SnS<sub>2</sub> structure has recently been reported to significantly enhance the PEC performance because of its higher specific surface area of catalytic sites and stronger light absorption by multiple photon scattering effect than a horizontal structure [9]. Additionally, the highly conductive edges of 2D SnS<sub>2</sub> nanosheets provide an efficient pathway for photoexcited carriers and good electronic contact with substrates. Recently, various synthesis methods, including chemical vapor deposition (CVD), hydrothermal reaction, and close space sublimation, have been utilized to synthesize SnS<sub>2</sub> nanosheets that vertically stand on substrates [9–15]. However, the methods have generally used solid powders as Sn and S precursors, such as stannic chloride (SnCl<sub>4</sub>·5H<sub>2</sub>O), SnS<sub>2</sub>, and S powders, which have high melting and boiling temperatures. As a result, 2D SnS<sub>2</sub> nanosheets have been formed at high growth temperatures near 450 °C, which considerably limit the conducting substrate choice. Essentially, further lowering the growth temperature will provide a useful path to exploit the benefits of various promising heterojunctions, such as vertically aligned 2D SnS<sub>2</sub> on graphene. Graphene has emerged

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as one of the most fascinating conductive layers, exhibiting a remarkable electron mobility ( $>15,000 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ) [16]. Liu et al. recently demonstrated the synthesis of vertically aligned  $\text{SnS}_2$  nanosheets on carbon cloth substrates for PEC water splitting [9].  $\text{SnS}_2$ /graphene nanocomposites have also been studied as a high rate anode material for lithium battery [17,18]. In addition, graphene and its derivatives have been extensively studied for various promising applications, including in-vivo bioimaging of  $\text{AgInZnS}$ -graphene oxide (GO) nanocomposites [19], X-ray photon detection of  $\text{CsPbBr}_3$ /reduced GO nanocomposites [20], and light-controlled conductive switching of  $\text{Cu}_2\text{O}$ /reduced GO nanocomposites [21]. Despite the considerable interest in and potential of 2D heterojunction systems, the growth and PEC performance of vertically aligned 2D  $\text{SnS}_2$  on graphene have not been demonstrated thus far. The graphene/ $\text{SnS}_2$  junction quality can crucially deteriorate owing to the high growth temperature of 2D  $\text{SnS}_2$  because graphene is just one or two layers of  $\text{sp}^2$ -bonded carbons. Thus, it is essential to grow 2D  $\text{SnS}_2$  at a low temperature in order to minimize the structural and chemical destruction of the graphene layer.

Additionally, the microscopic structure of 2D  $\text{SnS}_2$  nanosheets, such as the height and thickness, significantly affects the electrochemical OER and HER performance. Velicky et al. reported a strong dependence of charge transfer kinetics and PEC activity on the layer thickness of 2D  $\text{MoS}_2$  [22]. Bulk  $\text{MoS}_2$  exhibited faster charge transfer kinetics than monolayer  $\text{MoS}_2$  because of its stronger total light absorption, which strongly depends on the number of  $\text{MoS}_2$  layers. Moreover, bulk  $\text{MoS}_2$  also yielded a higher PEC photocurrent than few-layer and monolayer  $\text{MoS}_2$ . Although the comparison was based on horizontal 2D  $\text{MoS}_2$  structures on substrates, the results suggested that the structural optimization of 2D  $\text{SnS}_2$  nanosheets is central for exploiting the full potential of vertically aligned 2D  $\text{SnS}_2$  nanosheets for efficient PEC applications. However, little information is available on the structural optimization of vertically aligned 2D  $\text{SnS}_2$  for efficient carrier transfer and the corresponding PEC properties.

Herein, we report the significantly enhanced PEC performance of 2D  $\text{SnS}_2$  nanosheets that were vertically grown on graphene substrates via metalorganic CVD (MOCVD) at a low temperature of 230 °C. We successfully manipulated the microscopic structure of 2D  $\text{SnS}_2$  nanosheets, including the thickness and height, in a controllable and practical manner by varying the ratio of Sn and S precursors. We determined the optimal structure of 2D  $\text{SnS}_2$  nanosheets for efficient charge transfer and PEC performance via a systematic study on various microscopic structures. The PEC performance of the optimized 2D  $\text{SnS}_2$  nanosheets was further improved upon insertion of a graphene layer.

## 2. Experimental

Graphene was synthesized on Cu foils (Alfa Aesar) using inductively coupled plasma (ICP) CVD with  $\text{CH}_4$  and  $\text{H}_2$  gases at 950 °C for 5 min. The ICP power and growth pressure were fixed at 200 W and 1 Torr, respectively. The synthesized graphene on Cu was transferred to an indium tin oxide (ITO) glass substrate. The CVD growth and transfer procedures of graphene were described in further detail elsewhere [23]. 2D  $\text{SnS}_2$  was grown on ITO and ITO/graphene substrates at 230 °C at a pressure of 8 Torr for 15 min using a MOCVD system with  $\text{SnCl}_4$  and  $\text{H}_2\text{S}$  gas (5 vol.% in a balance of  $\text{N}_2$ ) as Sn and S precursors, respectively (Fig. 1).  $\text{SnCl}_4$  was vaporized at 7 °C and was then delivered into a quartz reaction tube using Ar gas. The flow rate ratios of the Ar carrier to  $\text{H}_2\text{S}$  gases were varied as 150:50, 125:75, 100:100, and 75:125, where the total flow rate of the gases was fixed at 200 SCCM (standard cubic centimeters per minute).

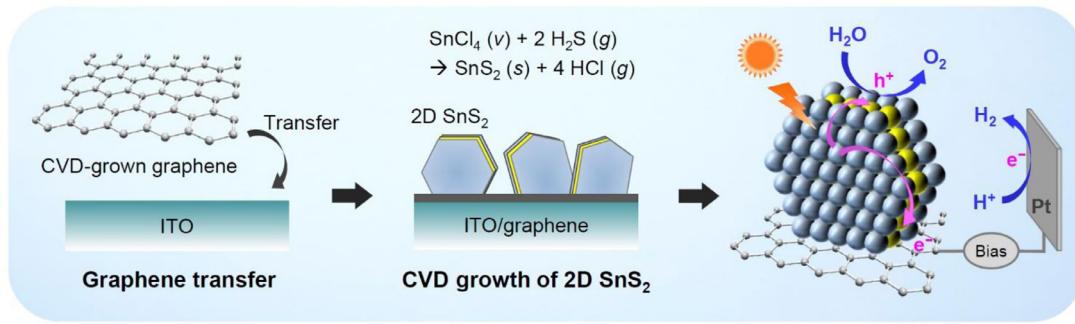
The 2D  $\text{SnS}_2$  morphology was characterized via scanning electron microscopy (SEM, Hitachi S-4800). The micro- and crystal structures of  $\text{SnS}_2$  were investigated using transmission electron microscopy (TEM, Tecnai G<sup>2</sup> F30 S-Twin), X-ray diffraction (XRD, Bruker AXS D8 Discover), and micro-Raman spectroscopy using an excitation band of 532 nm and a charge-coupled device detector. The optical properties were evaluated using an ultraviolet-visible (UV-vis) spectrophotometer (S-3100, SCINCO). The chemical states and sample compositions were characterized by X-ray photoelectron spectroscopy (XPS, Thermo Fisher K-Alpha+). The photoexcited carrier behavior was investigated by time-resolved photoluminescence (TRPL) measurements. The energy level of  $\text{SnS}_2$  was evaluated via UV photoelectron spectroscopy (UPS; Thermo scientific, K-alpha<sup>+</sup>).

PEC cells were fabricated on  $1 \times 2 \text{ cm}^2$  ITO and ITO/graphene substrates. The working area of the PEC photo-electrode was fixed at  $0.5 \times 0.5 \text{ cm}^2$  using nonconductive epoxy to cover the undesired area [24,25]. PEC characterization was performed using a three-electrode system and an electrochemical analyzer (potentiostat/galvanostat 263A). A Pt plate and KCl-saturated calomel ( $\text{Hg}/\text{Hg}_2\text{Cl}_2$ ) were used as the counter and reference electrodes, respectively. The electrolyte solution was 0.5 M  $\text{Na}_2\text{SO}_4$ . The light source used was a 150-W Xe arc lamp that delivered an intensity of  $100 \text{ mW/cm}^2$  of simulated AM 1.5 G irradiation. The current-voltage characteristics were recorded using a sourcemeter (Keithley 2400). Electrochemical impedance spectroscopy (EIS) measurements were performed under dark and constant light illumination ( $100 \text{ mW/cm}^2$ ) at a bias of 1.0 V while varying the AC frequency in the range from 100 kHz to 100 mHz. The incident monochromatic photon-to-current conversion efficiency (IPCE) of the electrode structure was measured using a grating monochromator in the excitation wavelength range of 300–700 nm.

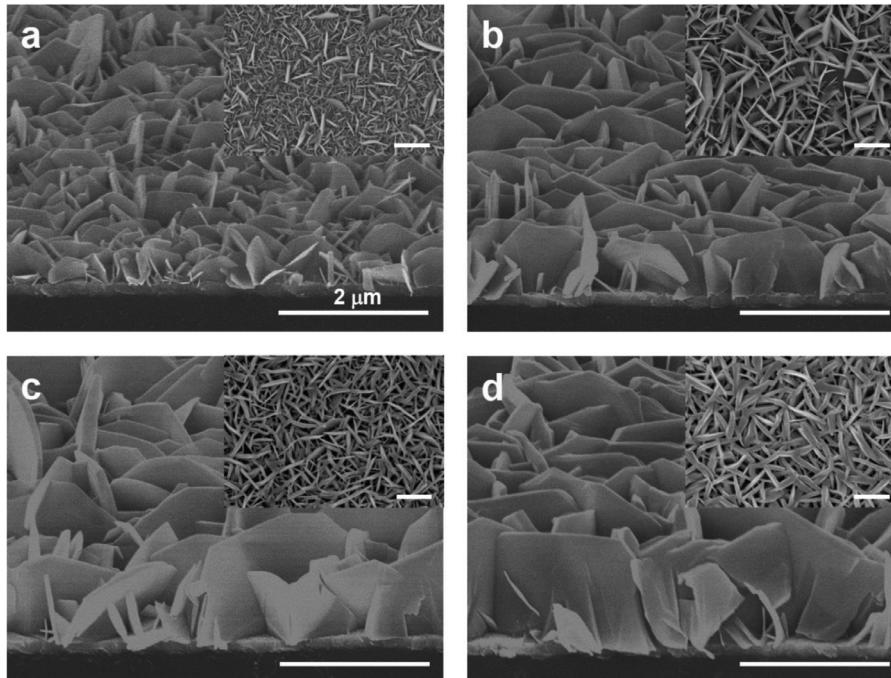
## 3. Results and discussion

The microscopic structure of 2D  $\text{SnS}_2$  was manipulated by the flow rate ratio of the Ar carrier gas to  $\text{H}_2\text{S}$  gas in a well-controlled manner (Figs. 2 and S1 in Supporting Information). The highest ratio of  $\text{SnCl}_4/\text{H}_2\text{S}$  (150:50 as the Ar carrier gas/ $\text{H}_2\text{S}$  flow rate ratio) produced the thinnest  $\text{SnS}_2$  nanosheets with an average thickness of  $\sim 25 \text{ nm}$  on an ITO substrate, which is hereinafter referred to as 2D  $\text{SnS}_2$ -S (Fig. 2a). The average height and width of the  $\text{SnS}_2$  nanosheets were  $\sim 0.5 \text{ } \mu\text{m}$  and  $\sim 1.0 \text{ } \mu\text{m}$ , respectively. Upon decreasing the ratio of  $\text{SnCl}_4/\text{H}_2\text{S}$ , the  $\text{SnS}_2$  nanosheets became larger and thicker. For a gas ratio of 125:75, the 2D  $\text{SnS}_2$  nanosheets appeared to be  $\sim 70 \text{ nm}$ ,  $\sim 0.85 \text{ } \mu\text{m}$ , and  $\sim 2.0 \text{ } \mu\text{m}$  in terms of thickness, height, and width, respectively, which is hereinafter referred to as 2D  $\text{SnS}_2$ -M (Fig. 2b). Upon further decreasing the gas ratio to 100:100, the 2D  $\text{SnS}_2$  nanosheets became larger,  $\sim 130 \text{ nm}$  in thickness,  $\sim 1.3 \text{ } \mu\text{m}$  in height, and  $\sim 2.5 \text{ } \mu\text{m}$  in width, which is hereinafter referred to as 2D  $\text{SnS}_2$ -L (Fig. 2c). For the gas ratio of 75:125, the thickness and height of 2D  $\text{SnS}_2$  nanosheets increased further to  $\sim 220 \text{ nm}$  and  $\sim 1.5 \text{ } \mu\text{m}$ , respectively, while the width remained almost unchanged at  $\sim 2.5 \text{ } \mu\text{m}$ , which is hereinafter referred to as 2D  $\text{SnS}_2$ -VL (Fig. 2d). The growth behavior implies that the  $\text{H}_2\text{S}$  gas may be a crucial facilitator for enhancing the chemical reaction, i.e.,  $\text{SnCl}_4 \text{ (v)} + 2 \text{ H}_2\text{S} \text{ (g)} \rightarrow \text{SnS}_2 \text{ (s)} + 4 \text{ HCl} \text{ (g)}$ , at such a low growth temperature [26]. The width of the 2D  $\text{SnS}_2$  nanosheets seemed to be self-constrained by being blockaded by other growing 2D nanosheets, as seen in 2D  $\text{SnS}_2$ -L sample. The height and thickness of 2D  $\text{SnS}_2$  nanosheets were then further increased at the gas ratio of 75:125 (2D  $\text{SnS}_2$ -VL).

The 2D  $\text{SnS}_2$  nanosheets showed typical XRD patterns of hexagonal  $\text{SnS}_2$  regardless of the gas ratios of  $\text{SnCl}_4/\text{H}_2\text{S}$  (Fig. 3a). The diffraction peaks at  $2\theta = 15.0^\circ, 28.2^\circ, 30.4^\circ, 32.1^\circ, 41.8^\circ, 45.6^\circ, 49.9^\circ$ , and  $52.4^\circ$  can be attributed to the (001), (100), (002), (011),



**Fig. 1.** Schematic of the preparation method for vertically aligned 2D SnS<sub>2</sub> nanosheets on graphene.

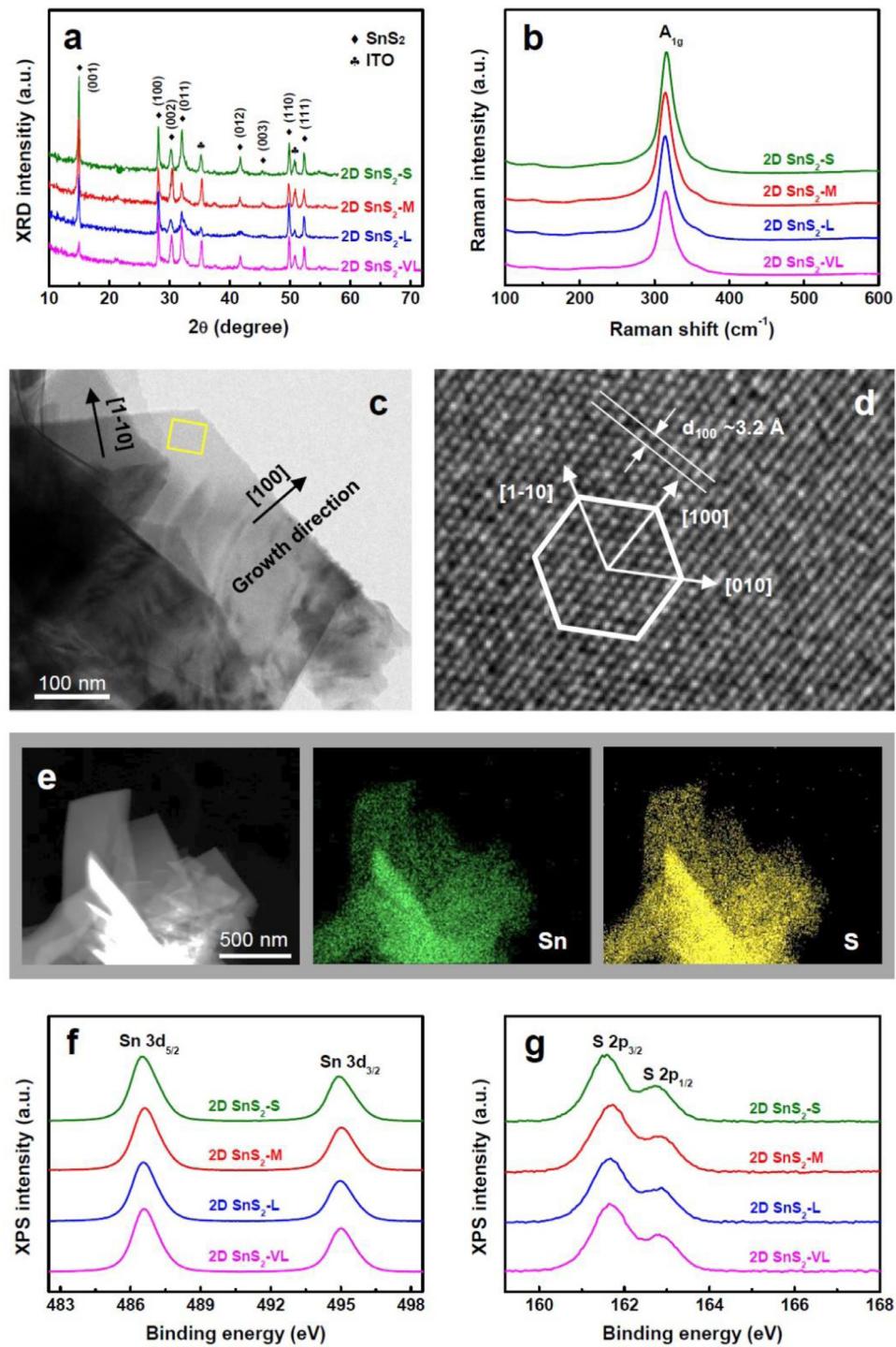


**Fig. 2.** SEM image of 2D SnS<sub>2</sub> nanosheets vertically grown on ITO substrates: (a) 2D SnS<sub>2</sub>-S, (b) 2D SnS<sub>2</sub>-M, (c) 2D SnS<sub>2</sub>-L, and (d) 2D SnS<sub>2</sub>-VL. The insets show the corresponding planar-view SEM images. Scale bar is 2  $\mu$ m.

(012), (003), (110), and (111) planes of the hexagonal structure of 2H-SnS<sub>2</sub> (JCPDS No. 89-2358). 2D SnS<sub>2</sub>-S and 2D SnS<sub>2</sub>-M exhibited the preferential (001) peak. With the increasing size of SnS<sub>2</sub> nanosheets (2D SnS<sub>2</sub>-L and 2D SnS<sub>2</sub>-VL), the (100) peak, corresponding to the vertical growth direction, became dominant. Meanwhile, all samples exhibited a strong Raman spectrum peak at 314 cm<sup>-1</sup>, which is attributed to the A<sub>1g</sub> mode, i.e., the vertical plane lattice vibration mode, of SnS<sub>2</sub> (Fig. 3b). The single peak confirmed the phase purity of SnS<sub>2</sub>. To further investigate the microstructure and crystallinity of the 2D SnS<sub>2</sub> nanosheets, a TEM study was performed. The TEM study revealed that most 2D SnS<sub>2</sub> nanosheets were single crystalline with sharp facets (Fig. 3c). The high-resolution TEM image, which was obtained from the corresponding yellow-boxed region in Fig. 3c, showed a lattice spacing of  $\sim$ 3.2 Å, corresponding to the (100) plane of SnS<sub>2</sub> (Fig. 3d) [27]. The lattice image exhibited a 6-fold rotation axis symmetry, indicating that the lattice image represented the (001) basal plane of the hexagonal SnS<sub>2</sub> structure. Thus, the mainly exposed surfaces of 2D SnS<sub>2</sub> nanosheets were (001) planes, which is consistent with the highest intensity of the (001) peak in the XRD patterns. However, the thin film-like morphology, formed by tightly-packed thick SnS<sub>2</sub> sheets, of 2D SnS<sub>2</sub>-VL exhibited a low relative intensity for the (001) peak. The energy-dispersive X-ray spectroscopy (EDX) el-

emental mapping clearly showed the uniform and homogeneous distribution of elemental Sn and S in the 2D nanosheets (Fig. 3e). An XPS study was performed to examine the chemical states of the 2D SnS<sub>2</sub> nanosheets. Notably, the 2D SnS<sub>2</sub> nanosheets exhibited the almost identical XPS spectra for the Sn 3d and S 2p core levels regardless of the gas ratio of SnCl<sub>4</sub>/H<sub>2</sub>S (Fig. 3f and g, respectively). The peaks at 486.6 eV and 495.0 eV are attributed to Sn 3d<sub>5/2</sub> and Sn 3d<sub>3/2</sub> of SnS<sub>2</sub>, respectively [28,29]. The single Gaussian peaks also indicated the phase purity of SnS<sub>2</sub>. In addition, the peaks at 161.6 eV and 162.8 eV corresponded to S 2P<sub>3/2</sub> and S 2P<sub>1/2</sub> of SnS<sub>2</sub>, respectively.

The PEC activities of the SnS<sub>2</sub> nanosheets were evaluated in 0.5 M Na<sub>2</sub>SO<sub>4</sub> solution by recording the linear sweep voltammograms under repetitive on-off switching of simulated AM 1.5G illumination (Fig. 4a). The photocurrent density kept increasing upon increasing the bias potential. However, the dark current became considerable above 1.3 V, which was attributed to the electrocatalysis and electrocorrosion in solution [24,30]. Interestingly, 2D SnS<sub>2</sub>-M exhibited the highest photocurrent density. Fig. 4b shows the apparent effect of SnS<sub>2</sub> microscopic size on PEC reactivity. 2D SnS<sub>2</sub>-M showed the highest photocurrent density of  $\sim$ 240  $\mu$ A/cm<sup>2</sup>, whereas 2D SnS<sub>2</sub>-L and 2D SnS<sub>2</sub>-VL, which possessed larger SnS<sub>2</sub> nanosheets, yielded significantly lower photocurrent densities of

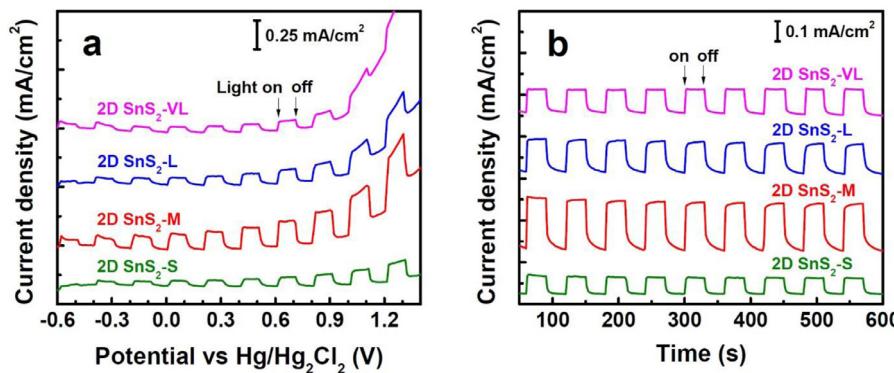


**Fig. 3.** (a) XRD patterns and (b) Raman spectra of 2D SnS<sub>2</sub>-S, 2D SnS<sub>2</sub>-M, 2D SnS<sub>2</sub>-L, and 2D SnS<sub>2</sub>-VL. (c) Low-magnification, (d) high-resolution lattice, and (e) EDX elemental mapping TEM images of 2D SnS<sub>2</sub>-M. XPS spectra of (f) Sn 3d and (g) S 2p core levels of 2D SnS<sub>2</sub>-S, 2D SnS<sub>2</sub>-M, 2D SnS<sub>2</sub>-L, and 2D SnS<sub>2</sub>-VL.

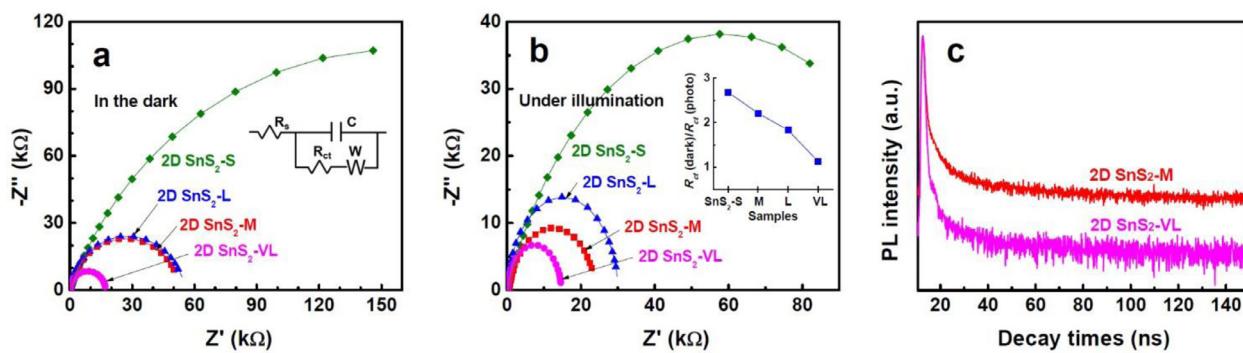
~150  $\mu\text{A}/\text{cm}^2$  and ~105  $\mu\text{A}/\text{cm}^2$ , respectively. 2D SnS<sub>2</sub>-S, which had the smallest and thinnest SnS<sub>2</sub> nanosheets, exhibited the lowest photocurrent density of ~90  $\mu\text{A}/\text{cm}^2$ .

To further understand the 2D SnS<sub>2</sub> nanosheet microscopic structure dependency of the PEC properties, systematic studies, including EIS, UV-vis absorption, and TRPL spectroscopy, were conducted. Fig. 5a and b show the EIS spectra of the 2D SnS<sub>2</sub> samples in the dark and under illumination, respectively. The Nyquist plots were fitted using a simplified Randles circuit (inset of Fig. 5a), which consists of the charge transfer resistance ( $R_{ct}$ ),

solution resistance ( $R_s$ ), constant phase element ( $Q$ ), and diffusion of species in the electrolyte solution represented by the Warburg impedance ( $W$ ). The  $R_{ct}$  values are provided in Table S1 in Supporting Information. 2D SnS<sub>2</sub>-S exhibited the largest radius of the EIS semicircle, indicating the highest charge transfer resistance in the dark ( $R_{ct}$  (dark) ~ 314 k $\Omega$ ), whereas 2D SnS<sub>2</sub>-VL showed the smallest radius ( $R_{ct}$  (dark) ~ 17 k $\Omega$ ). 2D SnS<sub>2</sub>-M and 2D SnS<sub>2</sub>-L possessed significantly lower  $R_{ct}$  (dark) values of ~53 k $\Omega$  and ~55 k $\Omega$ , respectively, than 2D SnS<sub>2</sub>-S. Under illumination, the EIS semicircle curvatures decreased owing to the contribution of



**Fig. 4.** (a) Photocurrent density–potential and (b) photocurrent density–time curves for PEC cells with various working electrodes (2D SnS<sub>2</sub>-S, 2D SnS<sub>2</sub>-M, 2D SnS<sub>2</sub>-L, and 2D SnS<sub>2</sub>-VL) in 0.5 M Na<sub>2</sub>SO<sub>4</sub> solution under repetitive on–off switching of simulated AM 1.5 G illumination.



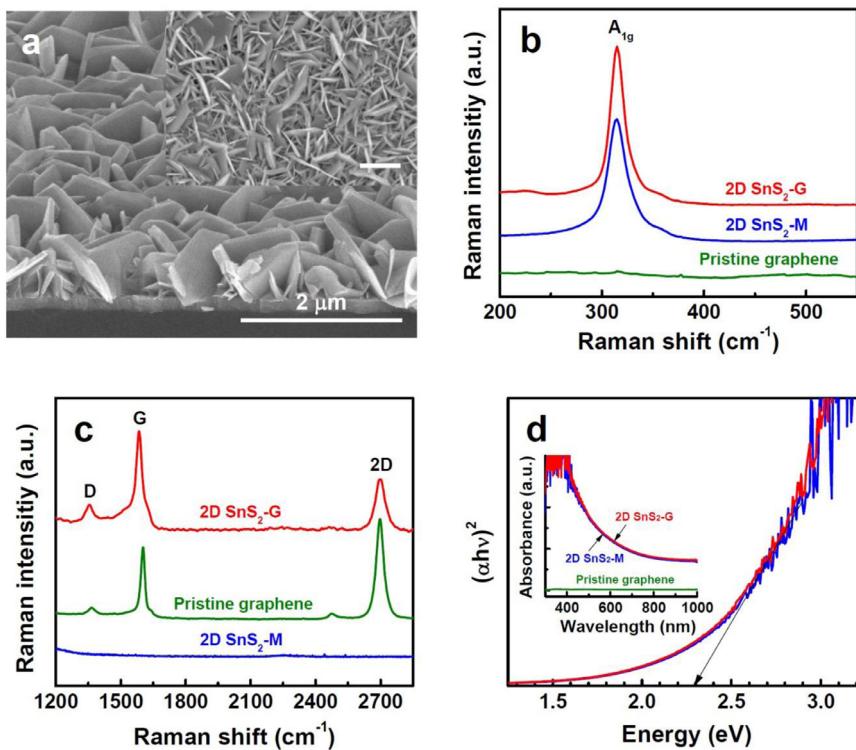
**Fig. 5.** Nyquist plots of 2D SnS<sub>2</sub>-S, 2D SnS<sub>2</sub>-M, 2D SnS<sub>2</sub>-L, and 2D SnS<sub>2</sub>-VL (a) in the dark and (b) under illumination. The inset in (a) shows an equivalent Randles circuit. (c) TRPL results of 2D SnS<sub>2</sub>-M and 2D SnS<sub>2</sub>-VL.

photogenerated charge carriers (Fig. 5b). Notably, the  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio continued decreasing with increasing size of the 2D SnS<sub>2</sub> nanosheets (inset of Fig. 5b). 2D SnS<sub>2</sub>-S still possessed a high  $R_{ct}$  (photo) value of  $\sim 117$  kΩ, despite having the highest  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio of 2.68. 2D SnS<sub>2</sub>-VL had the lowest  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio of 1.13, indicating the smallest gain by photogeneration of electron–hole (e–h) pairs. Of note, 2D SnS<sub>2</sub>-M had a lower  $R_{ct}$  (photo) value of  $\sim 24$  kΩ than 2D SnS<sub>2</sub>-L ( $\sim 30$  kΩ) owing to its higher  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio of 2.21.

The  $R_{ct}$  (dark) value difference between the samples is mainly determined by the ITO surface coverage of the 2D SnS<sub>2</sub> nanosheets, i.e., the specific interface area of ITO/SnS<sub>2</sub>. The larger specific interface area of thick 2D SnS<sub>2</sub> nanosheets results in a smaller  $R_{ct}$  value. Meanwhile, the  $R_{ct}$  (photo) values are crucially affected by the charge carrier dynamics, including the photogeneration rate of the e–h pairs upon light absorption and the transport efficiency of the carriers. The highest  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio of 2D SnS<sub>2</sub>-S indicates that photogenerated carriers efficiently transfer to the electrode with suppressed recombination. In the thinnest SnS<sub>2</sub> nanosheets, most e–h pairs are photogenerated near the surface and edges so that the e–h pairs can be readily separated from each other by the built-in surface potential. The carriers also have a shorter migration length to surface catalytic sites for PEC reaction [31]. Moreover, charge carriers are transported more efficiently to the electrode along the highly conductive surface and edges than through the bulk interior [32]. Nevertheless, the small specific interface area of 2D SnS<sub>2</sub>-S resulted in the highest  $R_{ct}$  (dark) and  $R_{ct}$  (photo) values, which were a decisive factor for the significantly poor PEC efficiency. In contrast, the bulk-like coverage of 2D SnS<sub>2</sub>-VL showed the lowest  $R_{ct}$  (dark) value. However, more e–h pairs not only are photogenerated inside the thick SnS<sub>2</sub> nanosheets than

near the surface but also migrate through the inefficient bulk interior. Thus, a poor PEC efficiency and the lowest  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio of 1.13 were observed. Obviously, the thickness of the SnS<sub>2</sub> nanosheets should be optimized by balancing a low interfacial contact resistance with effective photogenerated carrier dynamics. Of the samples, 2D SnS<sub>2</sub>-M with  $\sim 70$ -nm-thick SnS<sub>2</sub> nanosheets had a relatively low  $R_{ct}$  (dark) value of  $\sim 53$  kΩ with a high  $R_{ct}$  (dark)/ $R_{ct}$  (photo) ratio of 2.21, resulting in the highest PEC performance.

A TRPL spectroscopy study was conducted to understand the dynamic behavior of photogenerated carriers (Fig. 5c). The average carrier lifetimes were extracted by the PL decay kinetics fitted by a bi-exponential decay profile [33]. Compared to 2D SnS<sub>2</sub>-VL (0.75 ns), 2D SnS<sub>2</sub>-M showed a longer minority carrier lifetime of 1.7 ns. The slow PL decay, i.e., long carrier lifetime, can be attributed to suppressed e–h recombination through efficient separation and transportation of charge carriers. The estimated carrier lifetime of 0.75–1.7 ns is consistent with the reported value for single-crystalline SnS<sub>2</sub> flakes [11,34]. Assuming the intrinsic mobility of single-crystalline SnS<sub>2</sub> flakes ( $\sim 300$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>) [11], the carrier diffusion length was estimated to be  $\sim 0.7$ –1.1 μm. A significant amount of carriers, which are photogenerated in the upper region of SnS<sub>2</sub> nanosheets beyond the carrier diffusion length, cannot efficiently contribute to PEC reaction. Of the samples, 2D SnS<sub>2</sub>-M and 2D SnS<sub>2</sub>-L had the optimal height corresponding to nearly the limit of the carrier diffusion length. However, 2D SnS<sub>2</sub>-M exhibited significantly higher PEC reactivity than 2D SnS<sub>2</sub>-L because 2D SnS<sub>2</sub>-M had a thickness of  $\sim 70$  nm that effectively manifests nanosized effects, such as efficient separation and transportation of photogenerated e–h pairs with suppressed recombination.



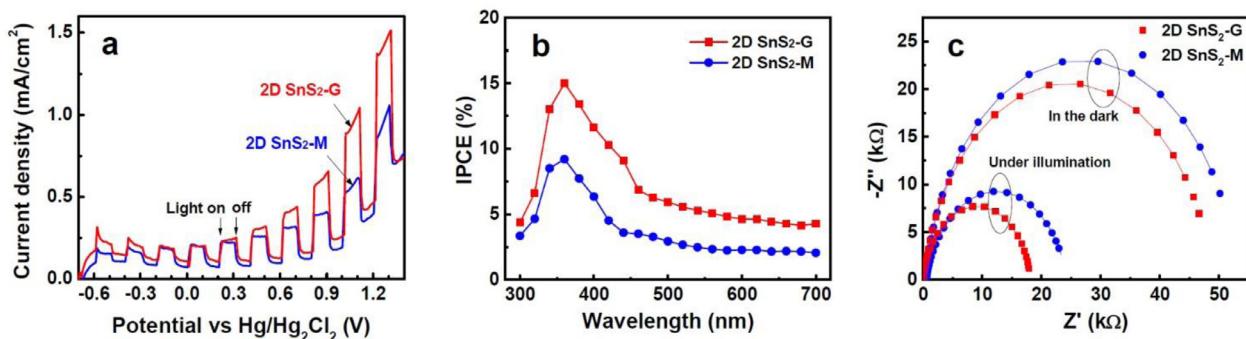
**Fig. 6.** (a) SEM image of 2D SnS<sub>2</sub> vertically grown on a graphene substrate (2D SnS<sub>2</sub>-G). Scale bar is 2  $\mu$ m. (b) and (c) Raman spectra of graphene, 2D SnS<sub>2</sub>-G, and 2D SnS<sub>2</sub>-M. (d) Plots of  $(\alpha h\nu)^2$  as a function of  $h\nu$  of 2D SnS<sub>2</sub>-G and 2D SnS<sub>2</sub>-M. The inset in (d) shows the corresponding UV-vis absorption spectra.

To further enhance the PEC performance of 2D SnS<sub>2</sub> nanosheets, they were grown vertically on graphene substrates under the same conditions as 2D SnS<sub>2</sub>-M. The 2D SnS<sub>2</sub> nanosheets on graphene, hereinafter referred to as 2D SnS<sub>2</sub>-G, showed nearly the same microscopic structure as 2D SnS<sub>2</sub>-M (Fig. 6a). The Raman spectrum of 2D SnS<sub>2</sub>-G also exhibited a single characteristic peak at 314 cm<sup>-1</sup>, which is almost identical to that of 2D SnS<sub>2</sub>-M (Fig. 6b). Owing to the low-temperature growth at 230 °C, the graphene layer was still present after MOCVD growth of SnS<sub>2</sub>, which was confirmed by the presence of characteristic G and 2D band peaks in the Raman spectrum (Fig. 6c). Compared to pristine graphene, 2D SnS<sub>2</sub>-G showed an increased  $I_D/I_G$  of ~0.17 and a decreased intensity ratio of 2D to G peaks owing to the surface coverage of SnS<sub>2</sub>. The CVD-grown graphene exhibited a low intensity ratio of characteristic D to G band peaks ( $I_D/I_G$  of ~0.1), indicating a high-quality graphene structure (Fig. 6c). The graphene showed excellent light transmittance of ~96.9% at 550 nm, corresponding to approximately one and a half layers of ideal graphene. Fig. 6d shows the UV-vis absorption spectra of 2D SnS<sub>2</sub>-M and 2D SnS<sub>2</sub>-G. Both samples showed almost identical light absorption characteristics, except for the slight light absorption difference from the graphene layer (~3%). The samples exhibited a bandgap energy of 2.30 eV, which was estimated from the intercept of the linear portion of the plot of  $(\alpha h\nu)^2$  as a function of energy ( $h\nu$ ).

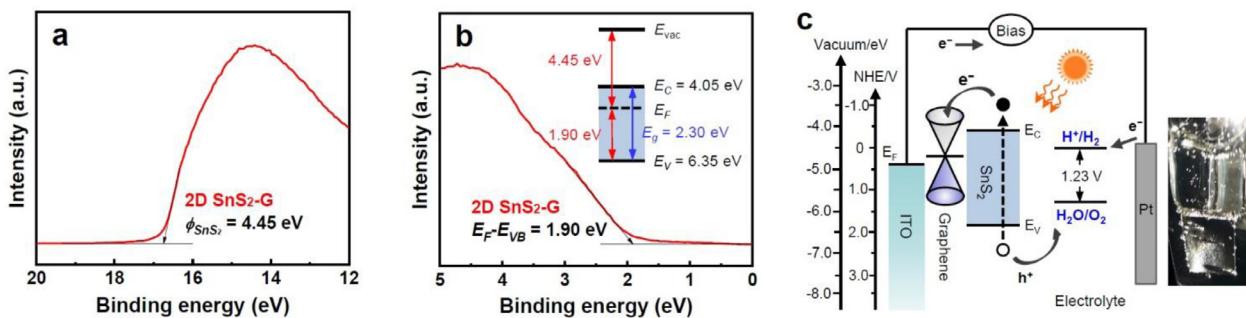
Fig. 7a shows the linear sweep voltammograms of 2D SnS<sub>2</sub>-M and 2D SnS<sub>2</sub>-G in 0.5 M Na<sub>2</sub>SO<sub>4</sub> solution under repetitive on-off switching of simulated AM 1.5G illumination. Both samples exhibited a significant photocurrent density even at zero bias potential. With increasing the bias potential, the photocurrent density of 2D SnS<sub>2</sub>-G increased more steeply than that of 2D SnS<sub>2</sub>-M. Eventually, 2D SnS<sub>2</sub>-G yielded approximately two times higher photocurrent density than 2D SnS<sub>2</sub>-M above 0.8 V. This considerably enhanced PEC activity was mainly attributed to the improved charge transfer properties upon insertion of the graphene layer. Fig. 7b shows the IPCE plots of two working electrodes (2D SnS<sub>2</sub>-M and

2D SnS<sub>2</sub>-G). 2D SnS<sub>2</sub>-G exhibited significantly improved photoconversion efficiencies over the overall incident-light waveband, suggesting fortified carrier photogeneration and transfer dynamics. The maximum photoconversion efficiencies of 9.2% and 15.0% were achieved at 360 nm for 2D SnS<sub>2</sub>-M and 2D SnS<sub>2</sub>-G, respectively. The photoconversion efficiency and photocurrent density of 2D SnS<sub>2</sub>-G were at least comparable with various SnS<sub>2</sub> photoanodes reported recently, such as SnS<sub>2</sub>/CuInS<sub>2</sub> nanosheet films [35], Al-doped ZnO inverse opal shell/SnS<sub>2</sub> core-shell structures [36] and SnO<sub>2</sub>/SnS<sub>2</sub> heterojunctions [37,38]. As shown in Fig. 7c and Table S1, the EIS Nyquist plot of 2D SnS<sub>2</sub>-G exhibited lower  $R_{ct}$  (dark) and  $R_{ct}$  (photo) values (49 k $\Omega$  and 18 k $\Omega$ , respectively) than those of 2D SnS<sub>2</sub>-M. The results indicate significantly improved charge transfer properties across the interface between SnS<sub>2</sub> nanosheets and ITO by inserting a graphene layer. The graphene layer not only acts as an excellent conducting layer [16] but also forms a favorable heterojunction with SnS<sub>2</sub> for efficient carrier transport. Notably, the  $R_{ct}$  (dark) to  $R_{ct}$  (photo) ratio of 2D SnS<sub>2</sub>-G was significantly improved to 2.72, suggesting that the graphene/SnS<sub>2</sub> heterojunction was more beneficial under illumination than in the dark.

To gain insight into the carrier transport property across the heterojunction of graphene/SnS<sub>2</sub>, its electronic structure was studied via UPS. The work function of SnS<sub>2</sub> ( $4.45 \pm 0.15$  eV) was determined on the basis of the difference between the photon energy of excited radiation (21.2 eV) and the spectrum width which is measured from the valence band and secondary edges (16.75 eV, Fig. 8a). The energy difference between the Fermi energy and valence band edge ( $E_F - E_{VB}$ ) was 1.90 eV (Fig. 8b). Given the bandgap energy of ~2.30 eV for SnS<sub>2</sub> based on the UV-Vis absorption spectrum, the electron affinity of SnS<sub>2</sub> was approximately 4.05 eV (inset in Fig. 8b). This electronic structure suggests the n-type behavior of SnS<sub>2</sub>, which works as a photoanode. As shown in Fig. 8c, the Fermi level (~4.6–4.8 eV of work function) of pristine few-layer graphene [39] was appropriately located between the Fermi level of the ITO and the conduction band edge of SnS<sub>2</sub> for efficient extraction of



**Fig. 7.** (a) Photocurrent density–potential for PEC cells with 2D SnS<sub>2</sub>-G and 2D SnS<sub>2</sub>-M as the working electrode in 0.5 M Na<sub>2</sub>SO<sub>4</sub> solution under repetitive on–off switching of simulated AM 1.5 G illumination. (b) IPCE plots and (c) Nyquist plots of the corresponding PEC cells.



**Fig. 8.** (a) UPS secondary electron cut-off and (b) valence spectra of 2D SnS<sub>2</sub>-G. The inset in (b) shows the corresponding energy band diagram of 2D SnS<sub>2</sub> nanosheets on graphene. (c) PEC water-splitting working principle of 2D SnS<sub>2</sub> nanosheets on graphene. The photo image in (c) shows gas bubbling on the dark cathode (Pt) for 2D SnS<sub>2</sub>-G during PEC measurement.

electrons to the cathode. A significant amount of gas bubbles was observed from both the photoanode (SnS<sub>2</sub>) and the dark cathode (Pt) during PEC measurement, suggesting that the photocurrent can be attributed to a water-splitting reaction (Fig. 8c).

#### 4. Conclusions

We demonstrated the vertical growth of 2D SnS<sub>2</sub> nanosheets on graphene substrates using MOCVD. Owing to the low temperature growth at 230 °C, the graphene layer was still present after SnS<sub>2</sub> growth and significantly reduced the charge transfer resistance. As a result, 2D SnS<sub>2</sub>-G showed approximately two times higher photocurrent density than 2D SnS<sub>2</sub>-M. In addition, the microscopic structure of the 2D SnS<sub>2</sub> nanosheets was successfully controlled by the ratio of SnCl<sub>4</sub>/H<sub>2</sub>S. The systematic microscopic structure dependency PEC studies revealed that the following sizes of 2D SnS<sub>2</sub> nanosheets were optimal: a height of ~0.85 μm, corresponding to nearly the limit of the carrier diffusion length, and a thickness of ~70 nm that not only exhibits a reasonably low  $R_{ct}$  (dark) value but also effectively manifests nanosized effects, such as efficient separation of photogenerated e–h pairs by the built-in surface potential and fast carrier transportation along the surface and edges of the 2D nanosheets. These results, including the vertical growth of 2D SnS<sub>2</sub> on graphene, provide a promising path for exploiting the full potential of 2D SnS<sub>2</sub> for various PEC applications.

#### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### CRediT authorship contribution statement

**Dong-Bum Seo:** Conceptualization, Methodology, Data curation, Software, Writing - original draft. **Tran Nam Trung:** Software. **Eui-Tae Kim:** Supervision, Conceptualization, Visualization, Data curation, Writing - review & editing.

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#### Supplementary materials

Supplementary material associated with this article can be found, in the online version, at doi:10.1016/j.electacta.2020.137164.

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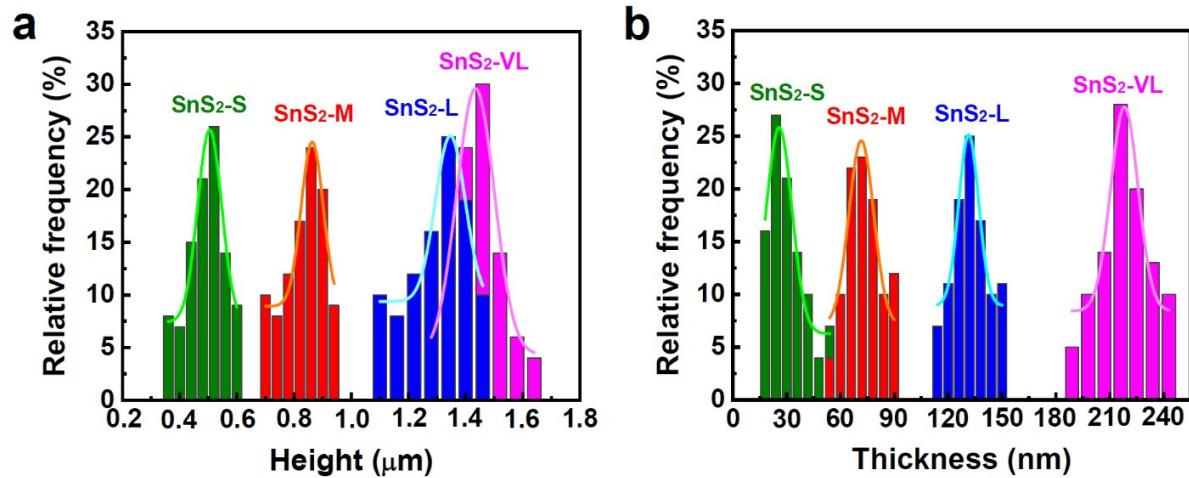
## Supporting Information for

### Controllable low-temperature growth and enhanced photoelectrochemical water splitting of vertical SnS<sub>2</sub> nanosheets on graphene

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**Fig. S1.** Height and thickness distributions of 2D SnS<sub>2</sub>-S, 2D SnS<sub>2</sub>-M, 2D SnS<sub>2</sub>-L, and 2D SnS<sub>2</sub>-VL.

**Table S1.**  $R_{ct}$  values for EIS analysis in the dark and under illumination.

Samples	$R_{ct}$ (dark) [kΩ]	$R_{ct}$ (photo) [kΩ]	$R_{ct}$ (dark)/ $R_{ct}$ (photo)
2D SnS <sub>2</sub> -S	314	117	2.68
2D SnS <sub>2</sub> -M	53	24	2.21
2D SnS <sub>2</sub> -L	55	30	1.83
2D SnS <sub>2</sub> -VL	17	15	1.13
2D SnS <sub>2</sub> -G	49	18	2.72